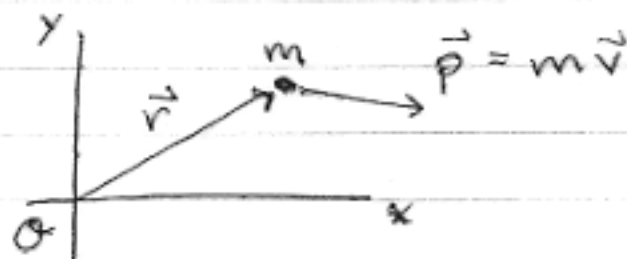


Angular Momentum (warm-up for H-atom)

Classically, angular momentum defined as (for a 1-particle system)

$$\vec{L} \equiv \vec{r} \times \vec{p}$$

$$= \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ x & y & z \\ p_x & p_y & p_z \end{vmatrix}$$



Note: \vec{L} defined w.r.t. an origin of coords

$$\vec{L} = \hat{x}(y p_z - z p_y) + \hat{y}(z p_x - x p_z) + \hat{z}(x p_y - y p_x)$$

(In QM, the operator corresponding to L_x is

$$\hat{L}_x = \hat{y} \hat{p}_z - \hat{z} \hat{p}_y, \quad p_z = \frac{\hbar}{i} \frac{\partial}{\partial z}, \text{ etc.}$$

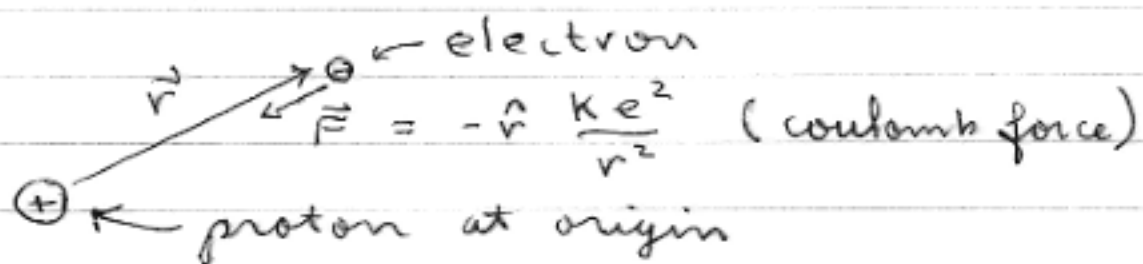
~ according to prescription of Postulate 2, part 3)

Classically, torque defined as $\vec{\tau} \equiv \vec{r} \times \vec{F}$, and

$$\vec{\tau} = \frac{d\vec{L}}{dt} \quad (\text{rotational version of } \vec{F} = m\vec{a})$$

If the force is radial (central force), then $\vec{\tau} = \vec{r} \times \vec{F} = 0 \Rightarrow \vec{L} = \text{const}$

H-atom:



In a multi-particle system, total ang. mom.

$$\vec{L}_{\text{tot}} = \sum_i \vec{L}_i \quad \text{is conserved for system isolated from external torques}$$

← sum over part's

Internal torques ^{can} cause exchange of ang. mom. among particles, but \vec{L}_{tot} remains constant.

In classical and quantum mechanics, only 4 things are conserved:

- energy
- linear momentum
- angular momentum
- electric charge

Back to QM. Define vector operator $\hat{\vec{L}}$

$$\hat{\vec{L}} = \hat{L}_x \hat{x} + \hat{L}_y \hat{y} + \hat{L}_z \hat{z} \quad \begin{array}{l} \leftarrow \text{operator} \\ \leftarrow \text{unit vector} \end{array}$$

Recall $\frac{d\langle Q \rangle}{dt} = \frac{i}{\hbar} \langle [\hat{H}, \hat{Q}] \rangle$

$$\Rightarrow \frac{d\langle \hat{\vec{L}} \rangle}{dt} = \frac{i}{\hbar} \langle [\hat{H}, \hat{\vec{L}}] \rangle$$

$$\frac{d\langle L_x \rangle}{dt} \hat{x} + \frac{d\langle L_y \rangle}{dt} \hat{y} + \dots$$

Claim: for a central force such as in H-atom

$$V = V(r) = -ke^2/r, \text{ then } [\hat{H}, \hat{\vec{L}}] = 0$$

(will show this later)

This implies $\frac{d\langle \hat{\vec{L}} \rangle}{dt} = 0$ (just like in classical mech.)

Angular momentum of electron in H-atom is constant, so long as it does not absorb or emit photon. Throughout present discussion, we ignore interaction of H-atom w/ photons.

Will show that for H-atom or for any atom, molecule, solid - any collection of atoms - the angular momentum is quantized in units of \hbar . $|\vec{L}|$ can only change by integer nbs of \hbar 's

$$\text{units of } L = [L] = [\hbar]$$

$$\text{Note } [L] = [r p], \quad [p] = \left[\frac{\hbar}{r}\right] \quad (\text{since } p = \hbar k)$$

$$\Rightarrow [L] = [r] \times [p] / [r] = [\hbar] \quad \checkmark$$

$$\text{Claim: } [\hat{L}_x, \hat{L}_y] = i\hbar \hat{L}_z$$

$$\text{and } \boxed{[\hat{L}_i, \hat{L}_j] = i\hbar \hat{L}_k}$$

(i, j, k cyclic:
 $x \ y \ z \text{ or}$
 $y \ z \ x \text{ or}$
 $z \ x \ y$)

To prove, need two very useful identities:

$$[A+B, C] = [A, C] + [B, C]$$

$$[AB, C] = A[B, C] + [A, C]B$$

$$\text{Proof: } [L_x, L_y] = [y p_z - z p_y, z p_x - x p_z] =$$

$$\underbrace{[y p_z, z p_x]}_{y [p_z, z] p_x} - \underbrace{[y p_z, x p_z]}_0 - \underbrace{[z p_y, z p_x]}_0 + \underbrace{[z p_y, x p_z]}_{x [z, p_z] p_y}$$

$\underbrace{[p_z, z]}_{-i\hbar}$
all other terms like $[y, p_x] = 0$
 $\underbrace{[z, p_z]}_{+i\hbar}$

$$= +i\hbar (x p_y - y p_x) = i\hbar L_z \quad \checkmark$$

(Have used $[x, p_x] = i\hbar$, $[x, y] = 0$, $[x, p_y] = 0$, $[p_x, p_y] = 0$, etc.)

I'm dropping the $\hat{}$ over operators when no danger of confusion.

Since $[L_x, L_y] \neq 0$, cannot have simultaneous eigenstates of \hat{L}_x and \hat{L}_y

$$\sigma_{L_x}^2 \sigma_{L_y}^2 \geq \left(\frac{1}{2i} \underbrace{\langle [\hat{L}_x, \hat{L}_y] \rangle}_{i\hbar \langle \hat{L}_z \rangle} \right)^2 = \left(\frac{\hbar}{2} \right)^2 \langle L_z \rangle^2$$

However, $L^2 = \vec{L} \cdot \vec{L} = L_x^2 + L_y^2 + L_z^2$ does commute w/ L_z

Claim: $[L^2, L_z] = 0$

$$\boxed{[L^2, L_i] = 0}, \quad i = x, y, \text{ or } z$$

$$\begin{aligned} \text{Proof: } [L^2, L_z] &= [L_x^2, L_z] + [L_y^2, L_z] + \underbrace{[L_z^2, L_z]}_0 \\ &= L_x \underbrace{[L_x, L_z]}_{-i\hbar L_y} + \underbrace{[L_x, L_z]}_{-i\hbar L_y} L_x + L_y \underbrace{[L_y, L_z]}_{+i\hbar L_x} + \underbrace{[L_y, L_z]}_{+i\hbar L_x} L_y \\ &= 0 \quad (\text{Note cancellations}) \end{aligned}$$

$[L^2, L_z] = 0 \Rightarrow$ can have simultaneous eigenstates of \hat{L}^2, \hat{L}_z (or \hat{L}^2, \hat{L}_i any i)

Looking forward to H-atom:

$$\hat{H} = -\frac{\hbar^2}{2m} \nabla^2(\cdot) + V(r) \cdot (\cdot)$$

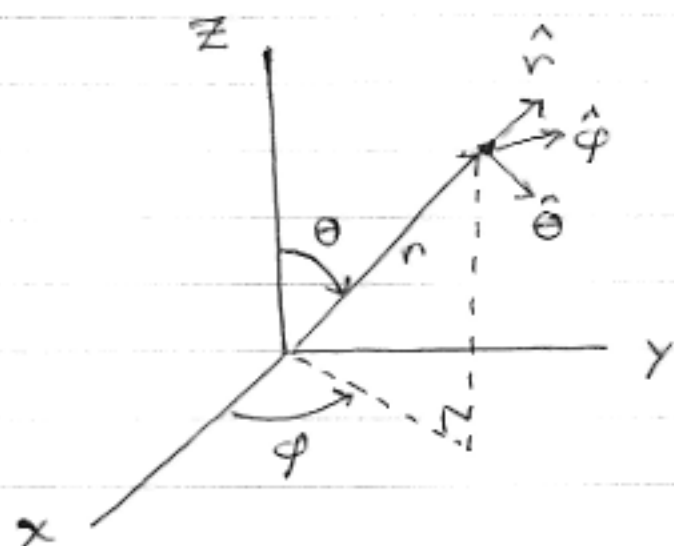
We will show that $[\hat{H}, \hat{L}^2] = 0, [\hat{H}, \hat{L}_z] = 0$

\Rightarrow simultaneous eigenstates of $\hat{H}, \hat{L}^2, \hat{L}_z$

$$\Psi = \Psi_{n \ell m}$$

\swarrow energy q-nbr
 \nwarrow L_z q-nbr
 \nwarrow L^2 q-nbr

When we solve the TISE $\hat{H}\Psi = E\Psi$ for the H-atom, the natural coords to use will be spherical coords: r, θ, φ (not x, y, z)



$$x = r \sin \theta \cos \varphi$$

$$y = r \sin \theta \sin \varphi$$

$$z = r \cos \theta$$

Just rewriting $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$ in

spherical coords is gawd-awful. But separation of variables will give special solutions, ~~esp~~ energy eigenstates, of form

$$\Psi(r, \theta, \varphi) = R(r) Y(\theta, \varphi) = R(r) \Theta(\theta) \Phi(\varphi)$$

The angular part of the solution $Y(\theta, \varphi)$ will turn out to be eigenstates of L^2 , L_z and will have form completely independent of the potential $V(r)$.

*

Given only $[L^2, L_z] = 0$ and \hat{L}^2, \hat{L}_z hermitean we know there must exist simultaneous eigenstates f (which will turn out to be the $Y(\theta, \varphi)$ mentioned above) such that

$$\hat{L}^2 f = \lambda \cdot f, \quad \hat{L}_z f = \mu \cdot f$$

(λ will be related to l , μ will be related to m)

We will show that f will depend on quantum-nbrs l, m , so we write it as f_l^m , and that

$$L^2 f_l^m = \hbar^2 l(l+1) \cdot f_l^m$$

$$L_z f_l^m = \hbar m \cdot f_l^m$$

where $l = 0, 1/2, 1, 3/2, \dots$ $m = -l, -l+1, \dots, l-1, l$

$f_l^m = Y_l^m(\theta, \phi)$ will be determined later

Notice max eigenvalue of $L_z (=l\hbar)$ is smaller than square root of eigenval of $L^2 = \hbar\sqrt{l(l+1)}$

So, in QM $L_z < |L|$ Odd!

Also notice $l=0, m=0$ state has zero angular momentum ($L^2=0, L_z=0$) so, unlike Bohr model, can have electron in state that is "just sitting there" rather than revolving about proton in H-atom