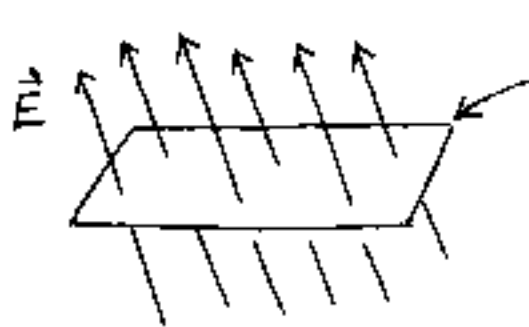


Knight Ch 27, Gauss's Law

G-1

New concept: electric flux Φ thru a surface

(begin w/ flat surfaces, uniform \vec{E})

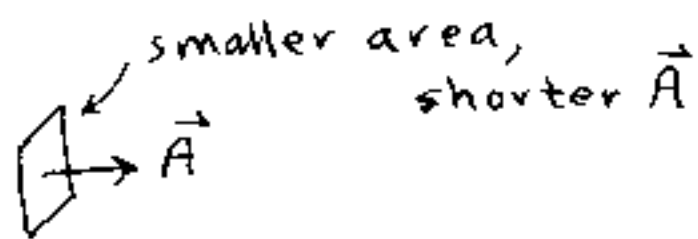
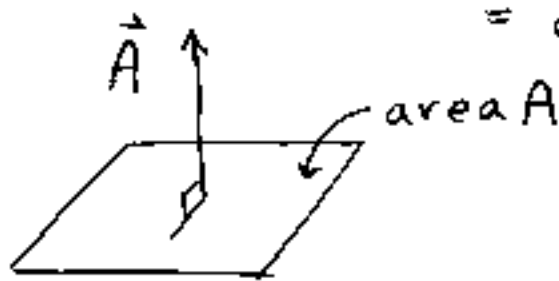



surface w/ area A has some electric flux thru it.

Define surface vector $\vec{A} = A \hat{n}$

Magnitude $A = |\vec{A}| = \text{area } A \text{ of surface}$

direction of \vec{A} = direction \perp to surface
= direction of unit normal \hat{n}

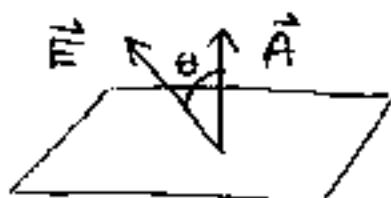


(Note there are 2 possible directions! )

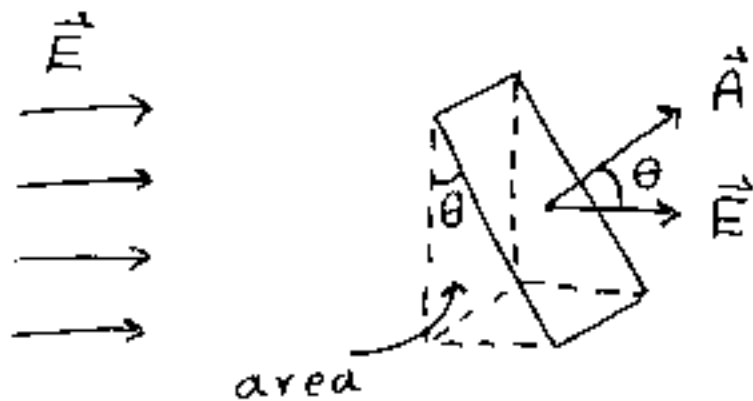
Flux Φ thru surface \vec{A} defined as

$$\Phi = \vec{E} \cdot \vec{A} = EA \cos \theta$$

← for special case $\vec{E} = \text{const}$, surface flat



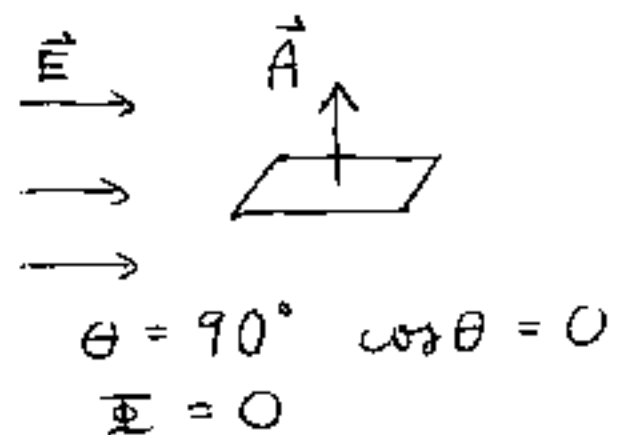
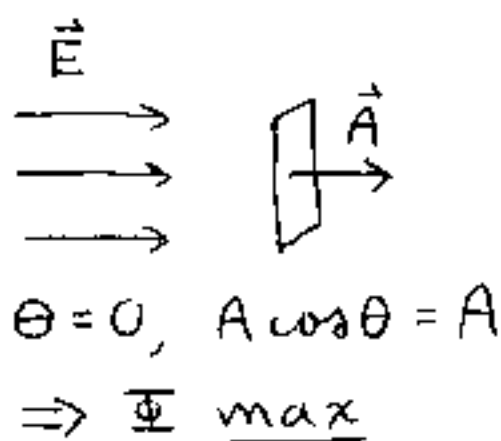
$$A \cos \theta = \text{area facing } \vec{E}$$



Φ large only if
 $A \cos \theta$ is large

$$A \cos \theta = \text{projection of area } \perp \text{ to } \vec{E}$$

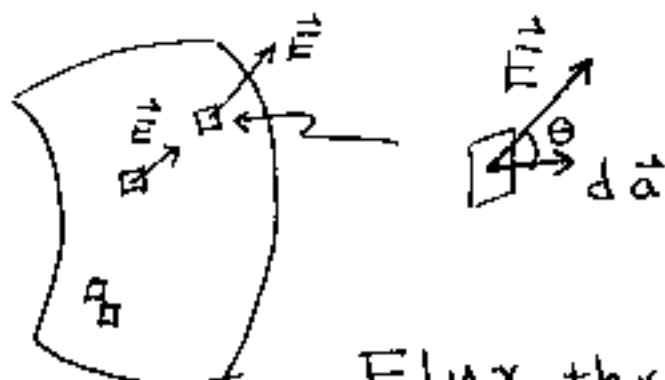
Think of field lines as rain flowing thru area \vec{A} . Flux = amount of rain thru area



Flux thru surface \propto # field lines crossing surface

If \vec{E} not const or area (surface) curved,

$$\Phi = \vec{E} \cdot \vec{A} \rightarrow \Phi = \int \vec{E} \cdot d\vec{a} = \text{"surface integral"}$$



Break surface up into many tiny segments of area da

$$\text{Flux thru segment } i = \Phi_i = \vec{E}_i \cdot d\vec{a}_i$$

Total flux thru whole surface =

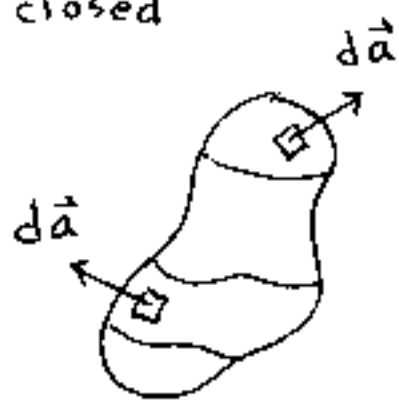
$$\Phi = \sum_i \Phi_i = \sum_i \vec{E}_i \cdot d\vec{a}_i \rightarrow \int \vec{E} \cdot d\vec{a}$$

Gauss's Law * (the 1st of 4 Maxwell's Eq'ns)

$$\oint \vec{E} \cdot d\vec{a} = \frac{Q_{in}}{\epsilon_0}$$

The electric flux thru any closed surface S is a constant ($1/\epsilon_0$) times the net charge Q enclosed by S .

$\oint \vec{E} \cdot d\vec{a}$ = "closed" surface integral
"closed"



← for closed surface,

direction of $d\vec{a}$ = outward
normal

ϵ_0 related to k by $k = \frac{1}{4\pi\epsilon_0}$

$\epsilon_0 = 8.85 \times 10^{-12}$ (SI units)

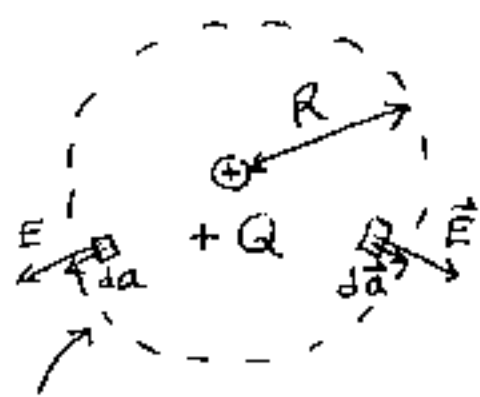
$$F_{coul} = \frac{k |q_1 q_2|}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{|q_1 q_2|}{r^2}$$

Gauss's Law can be derived from Coulomb's Law if charges are stationary, but Gauss's Law is more general than Coulomb's Law. Coulomb's Law is only true if q 's are stationary. Gauss's Law is always true whether or not q 's are moving.

Easy to show Gauss's Law is consistent w/ Coulomb's Law. Coulomb says...

$$E = \frac{kQ}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \quad (\text{for point charge } Q)$$

Does Gauss agree?



imaginary surface S
sphere, radius R

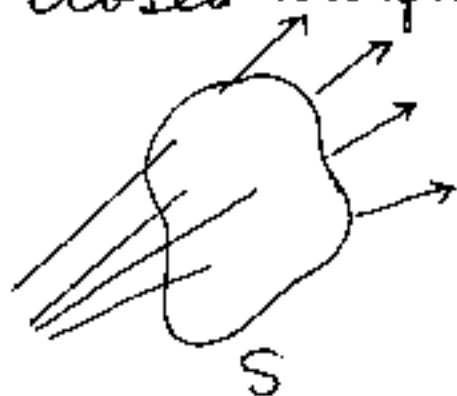
$$\begin{aligned} \oint_S \vec{E} \cdot d\vec{a} &= \int E da \\ & \quad (\text{since } \vec{E} \parallel d\vec{a}) \\ &= E \int da \quad (\text{since } E \text{ const}) \\ &= E \cdot A = E 4\pi R^2 \\ &= \frac{Q}{\epsilon_0} \quad (\text{says Gauss}) \end{aligned}$$

$$\Rightarrow E = \frac{Q}{4\pi\epsilon_0 R^2} \quad \checkmark$$

For closed surface ($d\vec{a}$ outward), flux is (+) over parts of surface where \vec{E} -field lines are exiting; flux is (-) where field lines are entering

The total (net) flux thru a closed surface $\Phi \propto [(\# \text{ lines exiting}) - (\# \text{ lines entering})]$

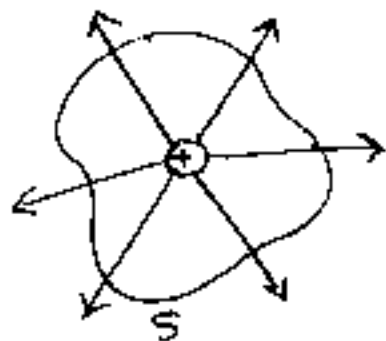
If field lines come from charge exterior to closed surface, then ($\# \text{ lines out} = \# \text{ lines in}$)



$$\Rightarrow \text{net flux} = 0$$

$$\oint_S \vec{E} \cdot d\vec{a} = 0$$

If field lines come from charge inside closed surface, then all field lines from charge exit surface

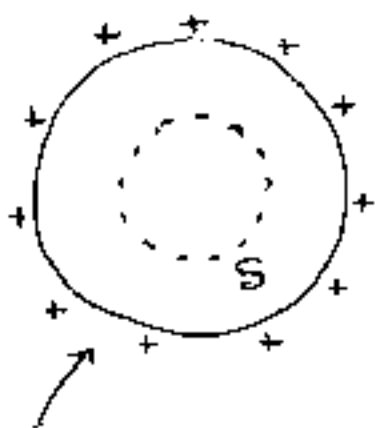


Total flux \propto
 $\#$ field lines originating inside
 \propto charge inside

$$\oint \vec{E} \cdot d\vec{a} = Q_{in} / \epsilon_0$$

In situations w/ high symmetry Gauss's Law allows quick calculation of \vec{E} -field.
Only easy to compute $\oint \vec{E} \cdot d\vec{a}$ if symmetric situation

Example: \vec{E} inside uniform, ^{spherical} shell of charge



uniform charge/area σ
on surface; no Q inside

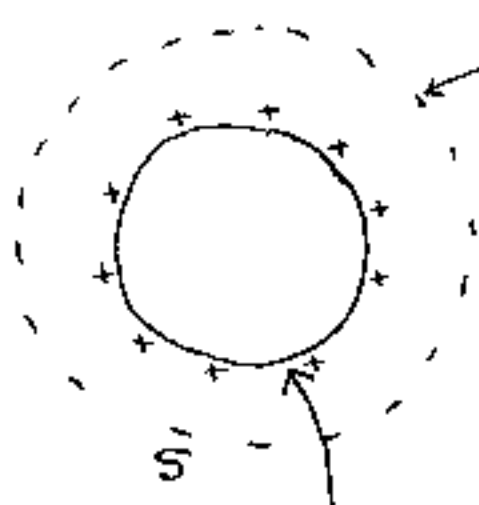
By symmetry, \vec{E} must
be radial (along a radius)
 $\Rightarrow \vec{E} = E(r) \hat{r}$

$$\Rightarrow \vec{E} \cdot d\vec{a} = \cancel{E da} E da$$

(since $\vec{E} \parallel d\vec{a}$)

$$\oint_S \vec{E} \cdot d\vec{a} = \oint E da = E \oint da = EA = \frac{Q_{in}}{\epsilon_0} = 0$$

$\Rightarrow \vec{E} = 0$ everywhere inside shell



total charge
 Q

With gaussian surface, ^S outside
shell, can show

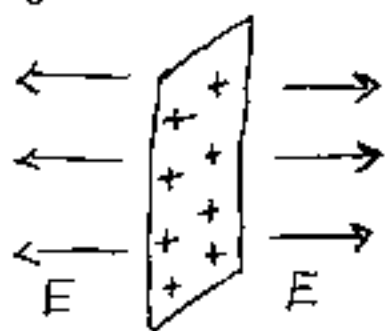
$$\vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r} \text{ everywhere}$$

outside
shell

\vec{E} -field near infinite, uniform plane of charge

$$\sigma = \frac{Q}{A} = \text{charge per area of plane}$$

"sigma" for "surface" (Knight uses η)



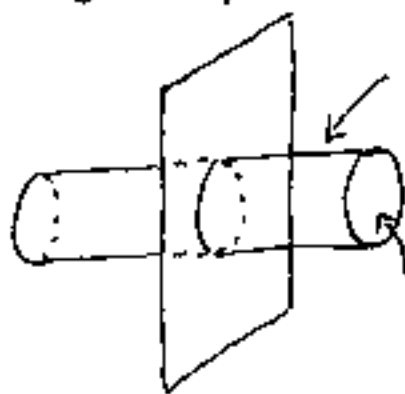
$$E = \frac{\sigma}{2\epsilon_0}$$

$\vec{E} \perp \text{plane}$

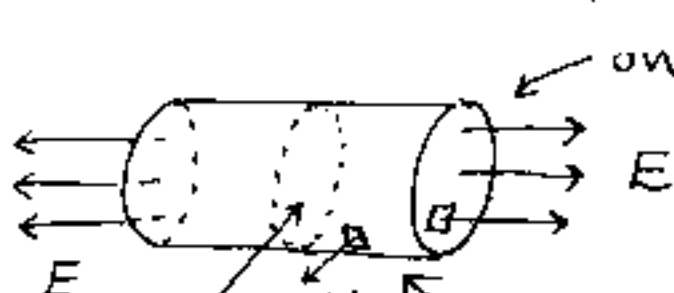
independent of distance from plane!

Proof: Apply Gauss's Law:

By symmetry, must have $\vec{E} \perp \text{plane}$



imaginary gaussian surface $S =$ cylinder
area A of endcap



on endcaps, $\vec{E} \cdot d\vec{a} = E da$
($\vec{E} \parallel d\vec{a}$)

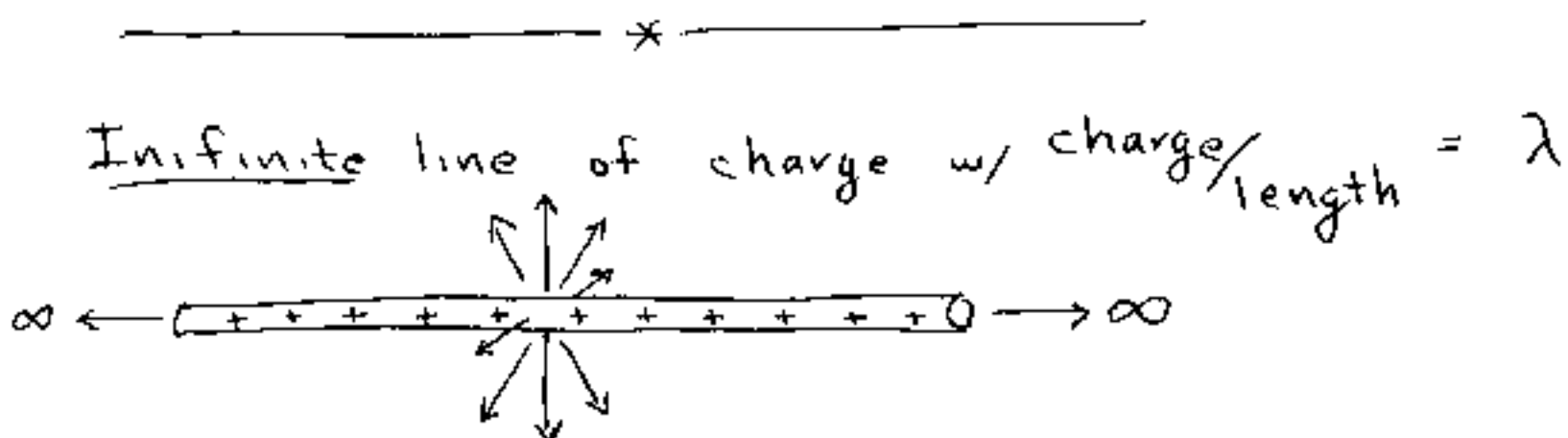
on sides, $\vec{E} \cdot d\vec{a} = 0$ ($\vec{E} \perp d\vec{a}$)

$$Q_{in} = \sigma \cdot A$$

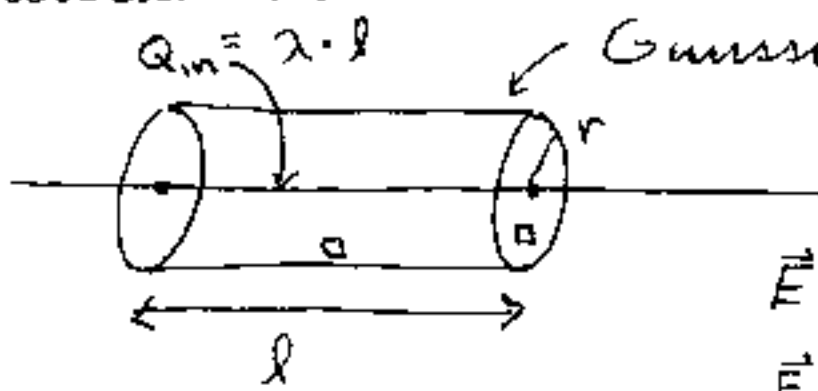
$$\oint_S \vec{E} \cdot d\vec{a} = \frac{Q_{in}}{\epsilon_0}$$

$$\oint_S = \int_{\text{left}} + \int_{\text{side}} + \int_{\text{right}} = EA + 0 + EA$$

$$\Rightarrow 2EA = \frac{\sigma A}{\epsilon_0} \Rightarrow E = \frac{\sigma}{2\epsilon_0} \quad \checkmark \quad \text{G-8}$$

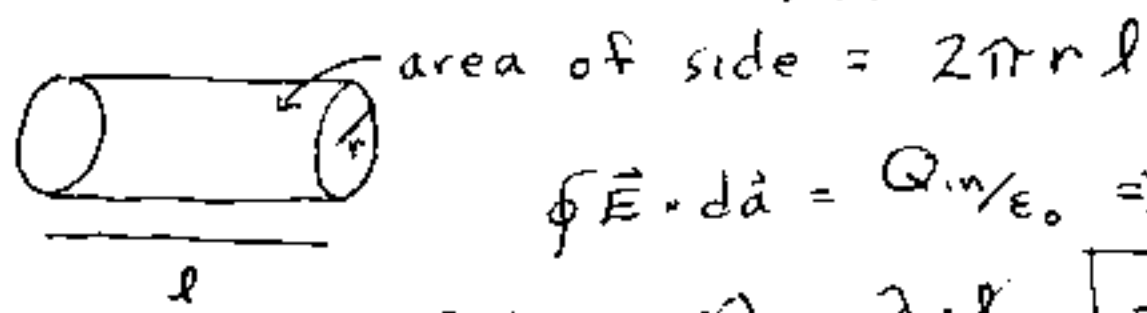


By symmetry, $\vec{E} \perp$ line $\Rightarrow E = E(r)$
 \vec{E} in radial direction, depends only on $r =$ radial distance



$\vec{E} \cdot d\vec{a} = 0$ on flat end caps
 $\vec{E} \cdot d\vec{a} = E da$ on curved side

$$\oint_S \vec{E} \cdot d\vec{a} = \int_{\text{curved side}} \vec{E} \cdot d\vec{a} = \int_{\text{curved side}} E da = E \int da = E A_{\text{side}}$$



$$\oint \vec{E} \cdot d\vec{a} = Q_{in}/\epsilon_0 \Rightarrow$$

$$E(2\pi r l) = \frac{\lambda \cdot l}{\epsilon_0}, \quad \boxed{E(r) = \frac{\lambda}{2\pi\epsilon_0 r}}$$

Conductors in electrostatic equilibrium.

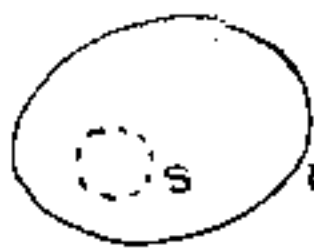
- $\vec{E} = 0$ inside a conductor (in equilibrium)
proof: If $\vec{E} \neq 0$ inside, then conduction e's feel a force $\vec{F} = q\vec{E}$ and will move.

Moving charge \Leftrightarrow not electrostat. equil.

In and on a conductor in equil., the charges arrange themselves so that $\vec{E} = 0$ everywhere inside (otherwise, not yet in equil.)

- The interior of a conductor in equil. can have no net charge (electrons + protons must have equal density).

proof from Gauss's Law: consider any closed surface S within conductor

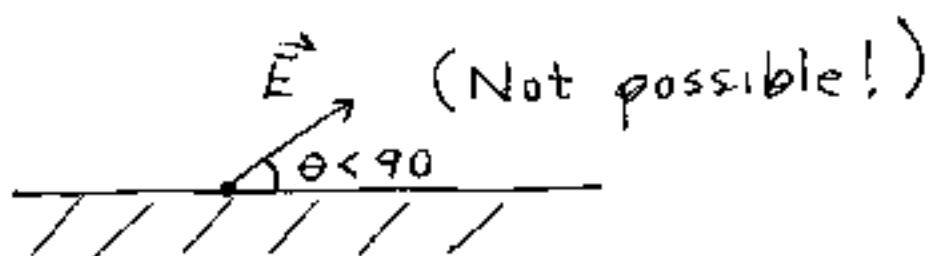
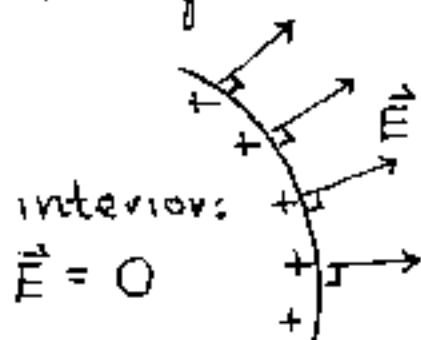


$$\oint_S \vec{E} \cdot d\vec{a} = 0 \quad (\text{since } \vec{E} = 0 \text{ on } S)$$

$$\Rightarrow Q_{in} = 0 \quad \checkmark$$

- If a conductor has a net charge Q , the charge must reside on the surface
- just showed it can't reside in interior \Rightarrow must be on surface
- \vec{E} -field near (at) surface of charged conductor must be \perp

\vec{E} -field near charged conductor must be \perp to surface.

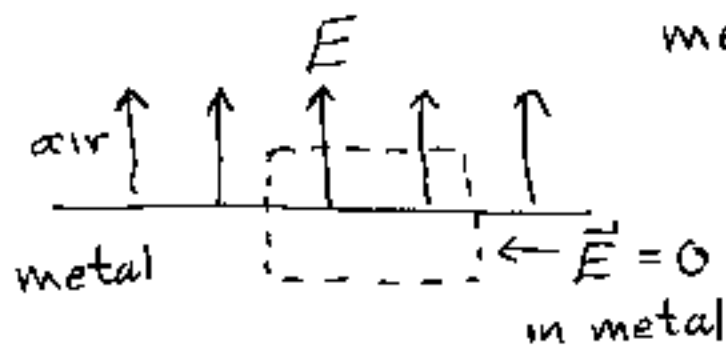
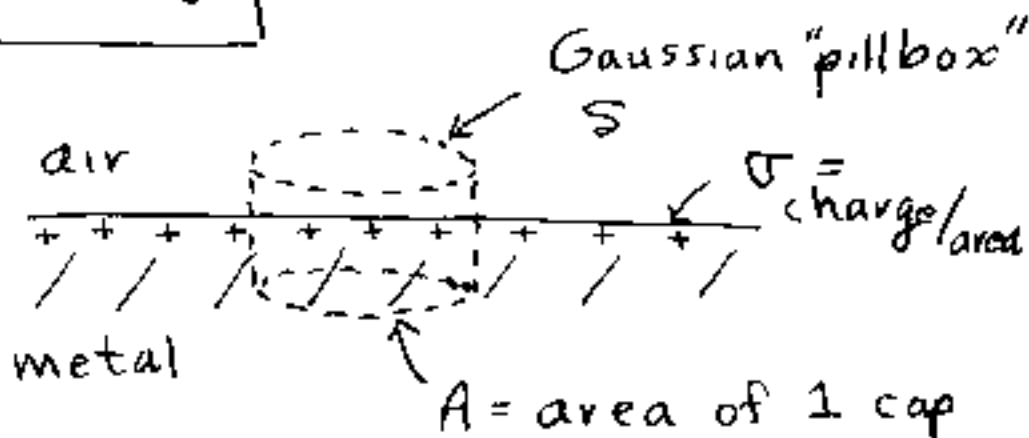


If \vec{E} not \perp to surface, there is a component of \vec{E} along surface ($E_x = E \cos \theta$), which would push conduction e's along surface
 $F_x = q E_x \Rightarrow$ would not have equilibrium.

- \vec{E} -field near surface of charged conductor has magnitude

$$E = \sigma / \epsilon_0$$

Proof by Gauss:



$$\oint_S \vec{E} \cdot d\vec{a} = \frac{Q_{in}}{\epsilon_0}$$

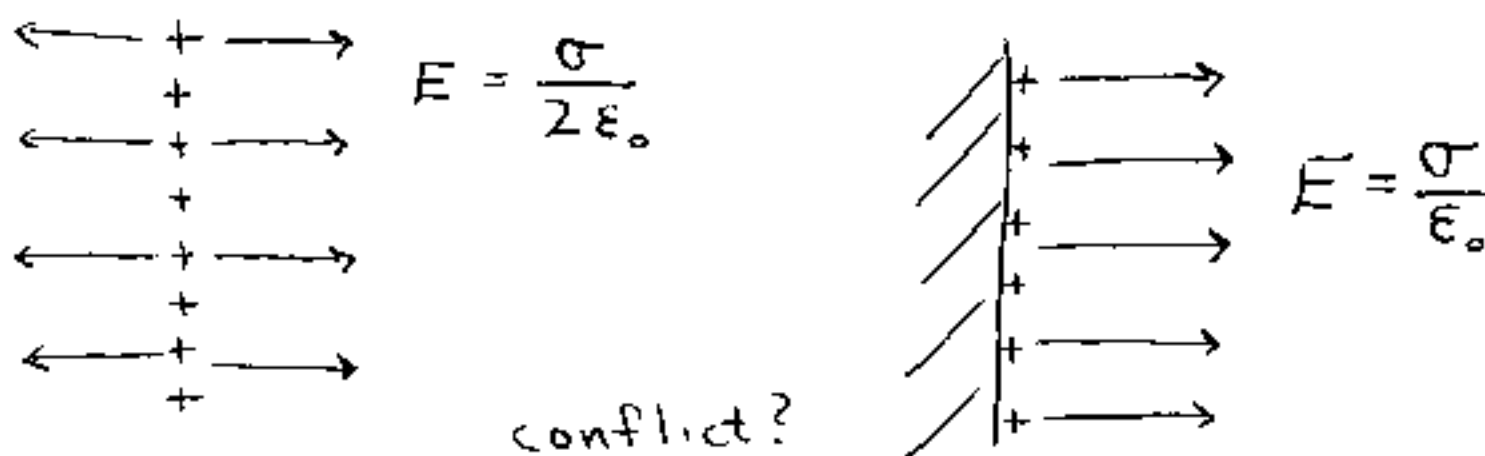
$$\Rightarrow E = \sigma / \epsilon_0$$

$$E A = \sigma A / \epsilon_0$$

(Notice similar, but different, from formula for \vec{E} due to plane of charge $\vec{E} = \frac{\sigma}{2\epsilon_0} \hat{n}$)

A puzzling situation has arisen!

Proved before that plane of charge creates field $E = \sigma / 2\epsilon_0$. Now have proved that plane of charge on surface of metal makes $E = \sigma / \epsilon_0$.



How can both be correct?

The \vec{E} -field anywhere is always due to all charges everywhere, including charges on far-away surfaces. Field near surface of metal is not only due to charges on nearby surface.

