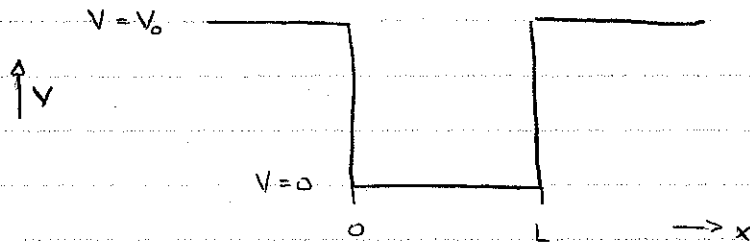


5.2 Finite square well

We will now extend our studies for a finite square well:



In the region $0 \leq x \leq L$ the solution is given (as before) by

$$\psi(x) = A \sin(kx) + B \cos(kx) \quad \text{with } k = \sqrt{\frac{m(E - 0)}{\hbar^2}}$$

Outside of this region we consider two different cases (note will not attempt to find the exact solution but we will discuss it qualitatively).

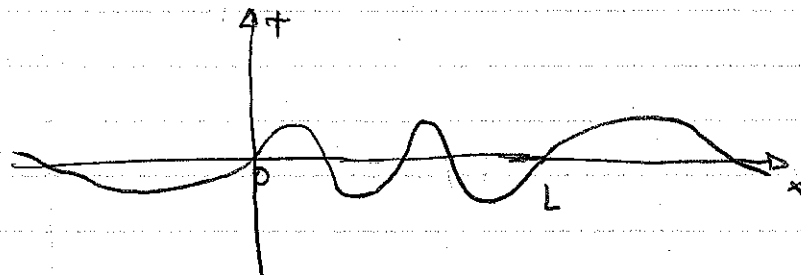
$E > V_0$: In this case the solution outside is also given by

$$\psi(x) = A_0 \sin(k_0 x) + B_0 \cos(k_0 x) \quad \text{with } k_0 = \sqrt{\frac{2m(E - V_0)}{\hbar^2}}$$

Thus, inside as well as outside we have oscillatory solutions, however with different wave numbers.

$$k_{\text{inside}} < k_{\text{outside}} \Rightarrow \lambda_{\text{inside}} > \lambda_{\text{outside}}$$

Qualitative sketch:



$E < V_0$ (but $E > 0$): Now the Schrödinger equation in the region outside $0 \leq x \leq L$ is given by

$$\frac{d^2 \psi}{dx^2} = +\beta^2 \psi(x) \quad \text{with} \quad \beta = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}}$$

positive ↗

Solutions are $\psi_1(x) = A_0 e^{\beta x}$ and $\psi_2 = B_0 e^{-\beta x}$

⇒ general solution: $\psi(x) = A_0 e^{\beta x} + B_0 e^{-\beta x}$ for $x < 0$

$$\psi(x) = A_0 e^{\beta(x-L)} + B_0 e^{-\beta(x-L)} \quad \text{for } x > L$$

Note that in the region

• $x < 0$: $e^{-\beta x}$ increases towards ∞ for $x \rightarrow -\infty$
and this is physically impossible

$$\Rightarrow \psi(x) = A_0 e^{\beta x} \quad \text{for } x < 0$$

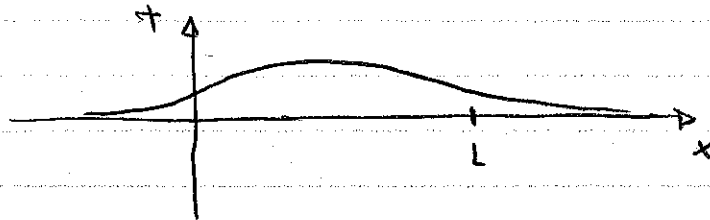
• $x > L$: $e^{\beta(x-L)}$ increases towards ∞ for $x \rightarrow \infty$
and this is physically impossible

$$\Rightarrow \psi(x) = B_0 e^{-\beta x} \quad \text{for } x > L$$

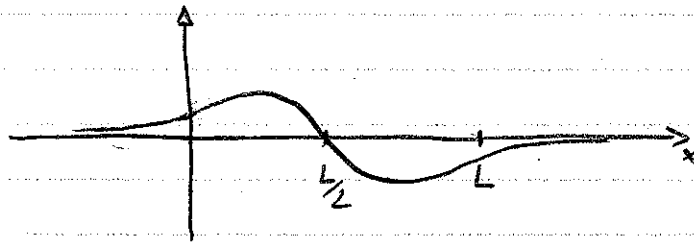
Thus, inside the well we have an oscillating solution (as before for the infinite square well), while outside it is an exponentially decaying solution.

Sketch of (stationary) wave functions for $E < V_0$.

$\psi_1(x)$:



$\psi_2(x)$:



Notes:

- For $E > V_0$ all energies are allowed.
For $E < V_0$ only discrete energies are allowed.
- Quantum mechanically, the particle can penetrate into the classically forbidden regions $x < 0$ and $x > L$.

One defines a characteristic distance = penetration depth as:

$$\frac{\psi(x = -\frac{1}{\beta})}{\psi(x=0)} = \frac{\psi(x=0) e^{+\beta(-\frac{1}{\beta})}}{\psi(x=0)} = e^{-1}$$

Thus, over the distance $\Delta x = \frac{1}{\beta}$ the wavefunction decays
by $e^{-1} = \frac{1}{e}$

$$\text{Note: } \Delta x = \frac{1}{\beta} = \frac{\sqrt{2m(V_0 - E)}}{\hbar}$$