## **Potentials**

10.1

Why can't we use a scalar potential to find the magnetic field, as we have done with the electric field, i.e., why can't we use

$$\mathbf{B} = -\nabla V_{B}(\mathbf{r})$$

- A. Because the divergence of **B** is always zero
- B. Because only either **E** or **B** can be described with a scalar potential, not both
- C. Because **B** can have a non-zero curl
- D. I don't know/remember
- E. None of the above (but I know the right reason!)

If I tell you  $\nabla \times \vec{F} = 0$ , what can you conclude about **F**?

- A)  $\vec{F} = 0$
- B)  $\vec{F} = \nabla f$  (for some f)
- C)  $\vec{F} = \nabla \cdot \vec{g}$  (for some  $\vec{g}$ )
- D)  $\vec{F} = \nabla \times \vec{g}$  (for some  $\vec{g}$ )
- E) Something else!

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- E) Something else!

If I tell you  $\nabla f = 0$ , what can you conclude about f?

- A) f = 0
- B)  $f = \nabla g$  (for some g)
- C)  $f = \nabla \cdot \vec{g}$  (for some  $\vec{g}$ )
- D)  $f = \nabla \times \vec{g}$  (for some  $\vec{g}$ )
- E) Something else!

$$\begin{split} \vec{B} &= \nabla \times \vec{A} \\ \vec{E} &= -\nabla V - \partial \vec{A} / \partial t \end{split}$$

$$\vec{B} = \nabla \times \vec{A}_{old}$$

$$\vec{E} = -\nabla V_{old} - \partial \vec{A}_{old} / \partial t$$

If I change gauge, so  $\vec{A}_{new} = \vec{A}_{old} + \nabla f$  I claim B is unaffected. But, what about E?

- A) Looks like E is also unaffected
- B) Looks like we changed E, but that's ok
- C) Looks like we changed E, that doesn't seem acceptable
- D) What are we doing?

$$\begin{split} \vec{A}_{new} &= \vec{A}_{old} + \nabla f \\ V_{new} &= V_{old} - \partial f / \partial t \end{split}$$

Gauge transformation

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla V - \partial \vec{A} / \partial t$$

10.1

Why can't we use a scalar potential to find the magnetic field, as we have done with the electric field, i.e., why can't we use

$$\mathbf{B} = -\nabla V_{R}(\mathbf{r})$$

- A. Because the divergence of **B** is always zero
- B. Because only either **E** or **B** can be described with a scalar potential, not both
- C. Because B can have a non-zero curl
- D. I don't know/remember
- E. None of the above (but I know the right reason!)

The Lorentz gauge is defined by:

$$\nabla \cdot \vec{A} = -\frac{1}{c^2} \frac{\partial V}{\partial t}$$

The main PDEs for the potentials are:

$$-\nabla^2 V + \frac{1}{c^2} \frac{\partial^2 V}{\partial t^2} = \frac{\rho}{\varepsilon_0}$$

$$-\nabla^2 \vec{A} + \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = \mu_0 \vec{J}$$

The solutions are:

$$V(\vec{r},t) = \frac{1}{4\pi\varepsilon_0} \iiint \frac{\rho(\vec{r}',t_R)}{|\vec{r}-\vec{r}'|} d^3\vec{r}' \text{ where } t_R = t - \frac{|\vec{r}-\vec{r}'|}{c}$$

$$\vec{A}(\vec{r},t) = \frac{\mu_0}{4\pi} \iiint \frac{\vec{J}(\vec{r}',t_R)}{|\vec{r}-\vec{r}'|} d^3 \vec{r}'$$

How have we been "setting the gauge" so far in this class?

- A) Requiring **A** and V go to 0 at infinity
- B) Putting a condition on **A** like: A=0, or  $\nabla \times A=0$
- C) Putting a condition on  $\nabla \cdot A$
- D) Choosing the function "f" in the formula  $A' = A + \nabla f$
- E) Something else, none of these, MORE than one, not really sure...

In Coulomb's gauge (CG): 
$$V(\vec{r},t) = \frac{1}{4\pi\varepsilon_0} \iiint \frac{\rho(\vec{r}',t)}{|\vec{r}-\vec{r}'|} d^3\vec{r}'$$

In Lorentz' gauge:

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$$V(\vec{r},t) = \frac{1}{4\pi\varepsilon_0} \iiint \frac{\rho(\vec{r}',t_R)}{|\vec{r}-\vec{r}'|} d^3\vec{r}' \qquad \text{where } t_R = t - \frac{|\vec{r}-\vec{r}'|}{c}$$

## These look subtly different, which is correct?

- A) CG is unphysical and incorrect, it violates relativity
- B) CG result is only correct for time independent problems, LG is what you want for time dependent problems
- C) They only LOOK different, but in fact they give the same result for V when you work them out
- D) Both are equally "correct", it's just a gauge choice. You can use either one in any situation
- E) Something else is going on, none of the above articulates my opinion very well here!

## How do you interpret

$$t_r \equiv t - \frac{\left| \vec{r} - \vec{r'} \right|}{c}$$

- A) is the actual time of observation at point r.
- B) is the time light needs to travel from r' to r.
- C) is a time in the future when light emitted from point r at time t arrives at point r'.
- D) is a time in the past such that light emitted from point r' arrives at r at time t.
- E) None of these.

## The 'retarded time',:

$$t_{r} \equiv t - \frac{\left| \vec{r} - \vec{r}' \right|}{c}$$

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- E) None of these.

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$$\vec{A}(\vec{r},t) = \frac{\mu_0}{4\pi} \iiint \frac{\vec{J}(\vec{r}',t_R)}{|\vec{r}-\vec{r}'|} d^3 \vec{r}'$$

At what time, t, does an observer at s first know the current was turned on?

- A) t=0

- B) t=c s
  C) t=s/c
  D) Other!!
  - E) Not sure about this?

At what time t, after s/c, does an observer at s see current from the entire wire?

- A) Immediately
- B) Never
- C) Something else

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- A) Immediately
- B) Never
  C) Something else

Only see contributions from regions inside the light travel distance:  $\sqrt{s^2 + z^2} \le ct$ 

Or 
$$|z| \le \sqrt{(ct)^2 - s^2}$$

$$I(t) = \begin{cases} 0 & \text{t} \le 0 \\ I_0 & \text{t} > 0 \end{cases}$$

$$\mathbf{A}(s,t) = \frac{\mu_0}{4\pi} \int_{-\infty}^{\infty} \frac{\mathbf{I}(z,t_R)}{\sqrt{s^2 + z^2}} dz \qquad t_R \equiv t - \frac{\sqrt{s^2 + z^2}}{c}$$

What is  $I(z,t_R)$ ?

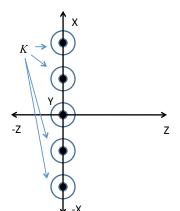
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$$\mathbf{A}(s,t) = \frac{\mu_0}{4\pi} \int_{-2?}^{2?} \frac{I_0 \hat{z}}{\sqrt{s^2 + z^2}} dz$$

What should the limits of integration be?

- A)  $-\infty$  to  $+\infty$
- B) -ct to +ct
- C)  $-\sqrt{c^2t^2+s^2}$  to  $+\sqrt{c^2t^2+s^2}$
- D)  $-\sqrt{c^2t^2-s^2}$  to  $+\sqrt{c^2t^2-s^2}$
- E) Something else, not sure, ...

A charge neutral and infinite static current sheet, **K**, flows in the x-y plane, in the y-axis direction. Therefore, to the right of the x-y plane, according to what you know from Phys 3310, the **E** and **B** field directions are:



- A) E along z-axis, B is zero
- B) **B** along *z*-axis, **E** is zero
- C) **B** along y, **E** along z
- D) **B** along *x*, **E** along *y*
- E) None of these

A charge neutral and infinite current sheet,  $\mathbf{K}$ , is turned on at t=0, flows in the x-y plane, in the y-axis direction. Therefore, immediately afterwards, the  $\mathbf{E}$  and  $\mathbf{B}$  fields:

- A) Remain zero (and for all time).
- B) Immediately appear with their static values in all space.
- C) Appear only near **K**
- D) Appear only to the right of **K**
- E) None of these

A charge neutral and infinite current sheet,  $\mathbf{K}$ , is turned on at t=0, flows in the x-y plane, in the y-axis direction. Therefore, shortly afterwards, the  $\mathbf{B}$  field near the sheet:

B=0 out there for awhile

I Front moves

- A) is in the z-direction
- B) is in the x-direction
- C) is in the y-direction
- D) is actually zero close to K.
- E) None of these

A charge neutral and infinite current sheet, **K**, is turned on at t=0, flows in the x-y plane, in the y-axis direction. Therefore, shortly afterwards, the E field

very near the sheet:

- A) is in the -z-direction
- B) is in the -x-direction
- C) is in the -y-direction
- D) is actually zero close to **K**.
- E) None of these

B=0out for

here awhile

A charge neutral and infinite current sheet, K, is turned on at *t*=0, flows in the *x-y* plane, in the *y*-axis direction. Therefore, shortly afterwards, the E field near the wavefront:

I Front moves at v.

- A) is in the -z direction B) is in the -x direction
  - C) is in the -y direction
  - D) is actually zero close to the front.
  - E) None of these

B=0 out here for awhile

